

Radiative X(3872) decays in a charmonium-molecule hybrid model

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X(3872) is the very interesting few body system !

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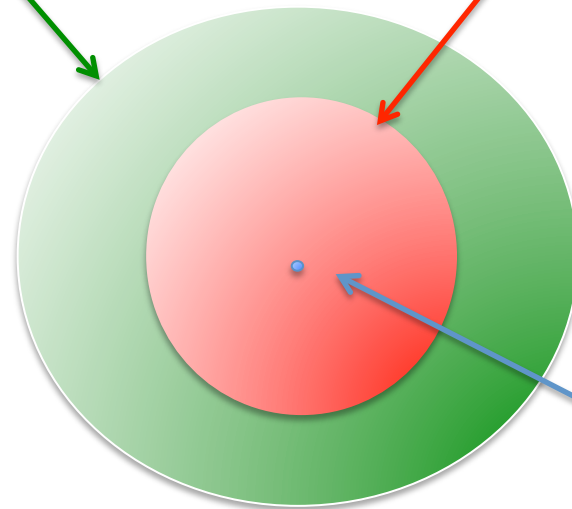
- Charmonium-molecule Hybrid model
- Why $X(3872)$ is so special?
- Radiative decays of $X(3872)$
- Radiative decay rates in our model
- Summary

Charmonium-molecule hybrid structure of X(3872)

J/psi rho and J/psi omega components are neglected for simplicity

$D^0 D^{*0\text{bar}}$

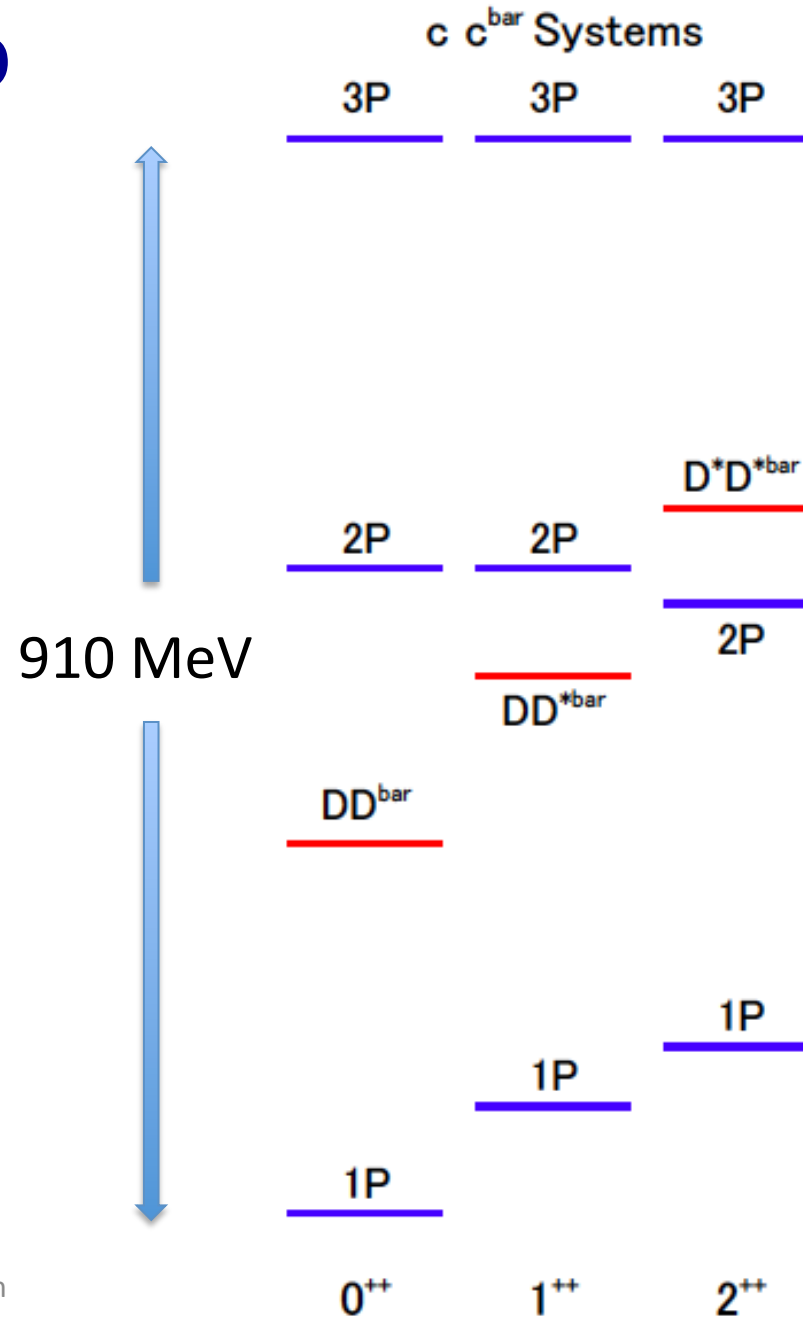
$D^+ D^{*-} + D^- D^{*+}$



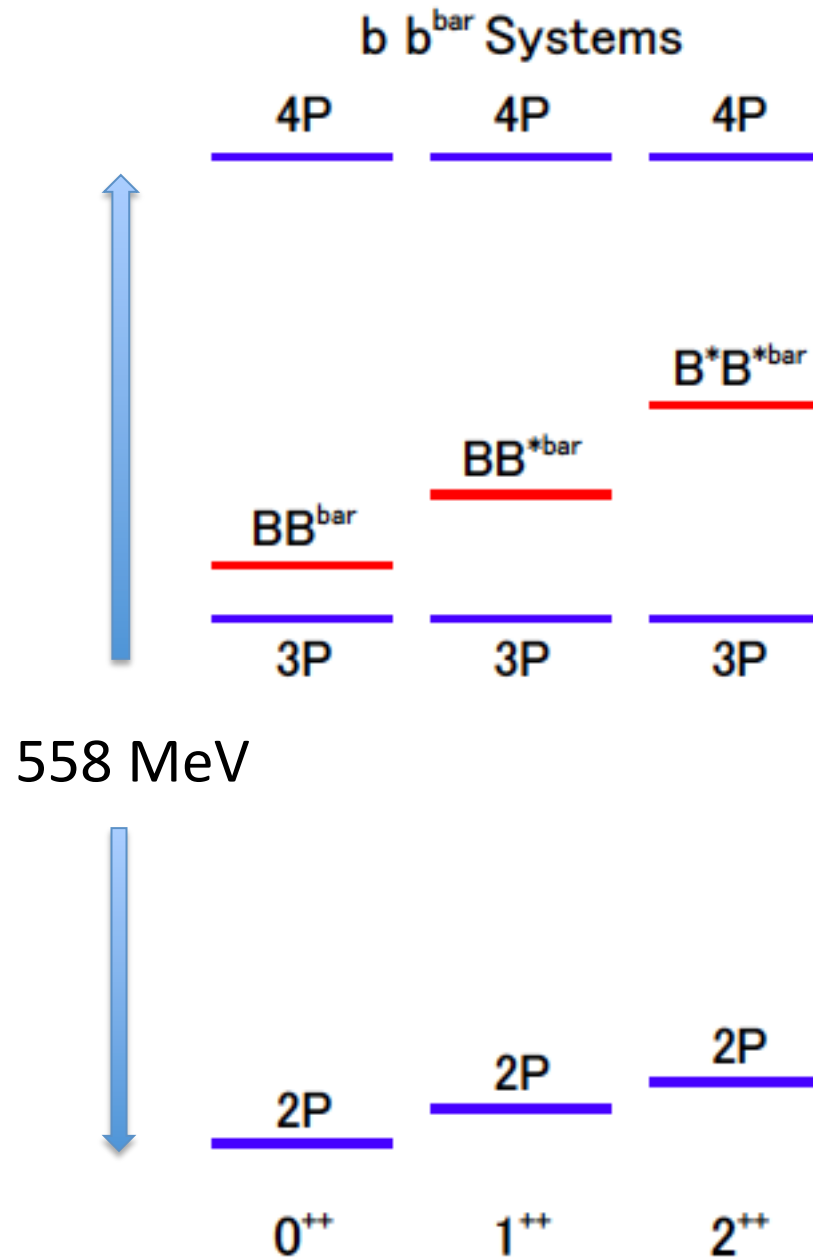
cc^{bar} charmonium

Ref. : Prog. Theor. Exp. Phys.
2013, 0903D01,

Why X(3872) is so special?



How about $b\bar{b}$ systems?



Radiative decays of $X(3872)$

- Experimental results

- Babar
$$R = \frac{B(X(3872) \rightarrow \psi' \gamma)}{B(X(3872) \rightarrow J / \psi \gamma)} = 3.4 \pm 1.4$$

Phys. Rev. Lett. 102, 132001 (2009)

- Belle
$$R = \frac{B(X(3872) \rightarrow \psi' \gamma)}{B(X(3872) \rightarrow J / \psi \gamma)} < 2.1$$

Phys. Rev. Lett. 107, 091803 (2011)

- Radiative decays are sensitive to the structure of $X(3872)$
- If $X(3872)$ is DD^* molecule, the transition probabilities of $X(3872) \rightarrow J/\psi \gamma, \psi' \gamma$ are very small
- If $X(3872)$ includes $J^{PC} = 1^{++}$ charmonium component, E1 transition is possible.
- $J^{PC} = 1^{++}$ charmonium is 2^3P_1 state: $\chi_{c1}(2P)$

- Matrix elements of E1 transitions are proportional to

$$\langle J/\psi | r | X(3872) : \chi_{c1}(2P) \rangle$$

$$\langle \psi' | r | X(3872) : \chi_{c1}(2P) \rangle$$

J/ψ is 1S and ψ' is 2S, so ψ' matrix element is larger -> support Babar's result

- If X(3872) has $\chi_{c1}(1P)$ component, situation will change, R becomes smaller

Coupling between C C-bar core, D⁰ D^{*0}-bar and D⁺ D^{*-}

- cc-bar core states: $|c\bar{c} : 1P\rangle$ $|c\bar{c} : 2P\rangle$
- D⁰ D^{*0}-bar state : $|D^0 \overline{D^{*0}}\rangle = \int d^3\vec{q} \varphi_0(\vec{q}) |D^0 \overline{D^{*0}}(\vec{q})\rangle$
- D⁺ D^{*-} state : $|D^+ D^{*-}\rangle = \int d^3\vec{q} \varphi_+(\vec{q}) |D^+ D^{*-}(\vec{q})\rangle$
in the center of mass frame
q is the conjugate momentum of the
relative coordinate

$$\langle D^0 \overline{D^{*0}}(\vec{q}') | D^0 \overline{D^{*0}}(\vec{q}) \rangle = \delta^3(\vec{q}' - \vec{q})$$

Coupling between C C-bar core, $D^0 \bar{D}^{*0}$ and $D^+ \bar{D}^{*-}$

- Charge conjugation + state is assumed
- Interaction: Isospin symmetric

$$\begin{aligned} \langle D^0 \bar{D}^{*0}(\vec{q}) | V | c\bar{c} : 1P \rangle &= \langle D^0 \bar{D}^{*0}(\vec{q}) | V | c\bar{c} : 2P \rangle \\ &= \frac{g \Lambda^2}{\vec{q}^2 + \Lambda^2} = \langle D^+ \bar{D}^{*-}(\vec{q}) | V | c\bar{c} : 1P \rangle = \langle D^+ \bar{D}^{*-}(\vec{q}) | V | c\bar{c} : 2P \rangle \end{aligned}$$

Coupling between C C-bar core, D⁰ D^{*0}-bar and D⁺ D^{*-}

- Schroedinger Equation

$$\begin{pmatrix} m_{c\bar{c}(1P)} - E & 0 & V & V \\ 0 & m_{c\bar{c}(2P)} - E & V & V \\ V & V & m_{D^0} + m_{D^{*0}} + \frac{\hat{p}^2}{2\mu_0} - E & 0 \\ V & V & 0 & m_{D^+} + m_{D^{*-}} + \frac{\hat{p}^2}{2\mu_+} - E \end{pmatrix} \begin{pmatrix} c_1 |c\bar{c}(1P)\rangle \\ c_2 |c\bar{c}(2P)\rangle \\ c_3 |D^0 \bar{D}^{*0}\rangle \\ c_4 |D^+ D^{*-}\rangle \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

$$\frac{1}{\mu_0} = \frac{1}{m_{D^0}} + \frac{1}{m_{D^{*0}}}, \quad \frac{1}{\mu_+} = \frac{1}{m_{D^+}} + \frac{1}{m_{D^{*-}}}$$

Numerical results: Mass

- Mass of the cc-bar core:
1P: 3.51 GeV 2P: 3.95 GeV
from S. Godfrey, N. Isgur, Phys. Rev. D 32 (1985)
189.
- Cutoff Lambda = 0.5 GeV,
Calculated bound states energy are
3.492 GeV and 3.87157 GeV
- $\chi_{c1}(1P)$ mass shift by coupling to DD^* is
about 20MeV

Numerical results: Wavefunction

$$\begin{aligned} |X(3872)\rangle = & 0.071|c\bar{c}(1P)\rangle - 0.326|c\bar{c}(2P)\rangle \\ & + 0.907|D^0\bar{D}^{*0}\rangle + 0.255|D^+D^{*-}\rangle \end{aligned}$$

- $c\bar{c}$ **1P** component is small.

E1 transition

- We have used results of the non-relativistic quark potential model by T. Barnes, S. Godfrey, E.S. Swanson, Phys. Rev. D72, 054026 (2005)
- Color Coulomb + linear scalar confinement + Gaussian smeared contact hyperfine interaction

- E1 transition width

$$\Gamma(X(3872) \rightarrow J/\psi(\psi') + \gamma) = \left(\frac{2}{3}\right)^4 \alpha \left| \langle \psi_f | r | \psi_i \rangle \right|^2 E_\gamma^3 \frac{E_f^{c\bar{c}}}{M_i^{c\bar{c}}}$$

$$\left| \langle J/\psi | r | \chi_{c1}(1P) \rangle \right|^2 = 4.13 \times 10^{-6}$$

$$\left| \langle J/\psi | r | \chi_{c1}(2P) \rangle \right|^2 = 1.49 \times 10^{-7}$$

$$\left| \langle \psi' | r | \chi_{c1}(1P) \rangle \right|^2 = 7.85 \times 10^{-6}$$

$$\left| \langle \psi' | r | \chi_{c1}(2P) \rangle \right|^2 = 1.08 \times 10^{-5}$$

$$\Gamma(X(3872) \rightarrow J/\psi + \gamma) = 29.2 \text{ keV}$$

$$\Gamma(X(3872) \rightarrow \psi' + \gamma) = 13.4 \text{ keV}$$

$$R = \frac{\Gamma(X(3872) \rightarrow \psi' + \gamma)}{\Gamma(X(3872) \rightarrow J/\psi + \gamma)} = 0.458$$

- Without $cc^{\text{bar}}:1$ 3P_1 component

$$|X(3872)\rangle = -0.293|c\bar{c}(2P)\rangle + 0.921|D^0\bar{D}^{*0}\rangle + 0.259|D^+D^{*-}\rangle$$

$$\Gamma(X(3872) \rightarrow J/\psi + \gamma) = 5.1 \text{ keV}$$

$$\Gamma(X(3872) \rightarrow \psi' + \gamma) = 7.7 \text{ keV}$$

$$R = \frac{\Gamma(X(3872) \rightarrow \psi' + \gamma)}{\Gamma(X(3872) \rightarrow J/\psi + \gamma)} = 1.50$$

Summary and outlook

- We have studied radiative decays of $X(3872)$.
- The calculated size of the $cc^{\text{bar}}:1\ ^3P_1$ component in $X(3872)$ is about 0.5%.
- Our result is consistent with experimental results.
- The ratio of the radiative decay widths is very sensitive to the size of the $cc^{\text{bar}}:1\ ^3P_1$ component in $X(3872)$.
- Effects of the J/ψ ρ and J/ψ ω components in $X(3872)$ should be studied.

Back up

About X(3872)

- First observation: 2003, Belle, KEKB cited more than 700 times
- $B^- \rightarrow K^- \pi^+ \pi^- J/\psi$ decay
Sharp peak of the invariant mass distribution of $\pi^+ \pi^- J/\psi$
- Mass: (3871.68 ± 0.17) MeV
about 0.2 MeV below $D^0 \bar{D}^{*0}$ -bar threshold
- Width: less than 1.2 MeV
- Quantum Number: $J^{PC} = 1^{++}$
- Other decay mode:
 $X(3872) \rightarrow \gamma J/\psi, \gamma \psi(2S)$
 $X(3872) \rightarrow \pi^+ \pi^- \pi^0 J/\psi$
 $X(3872) \rightarrow D^0 \underline{D}^{*0}$

Charmonium-molecule hybrid model of X(3872)

- Our works:
M. Takizawa and S. Takeuchi, Prog. Theor. Exp. Phys. 2013, 0903D01, **Hadronic model approach**
- Other groups works:
R. D. Matheus, F. S. Navarra, M. Nielsen, and C. M. Zanetti, Phys. Rev. D, 80, 056002 (2009). **QCD sum rule approach**

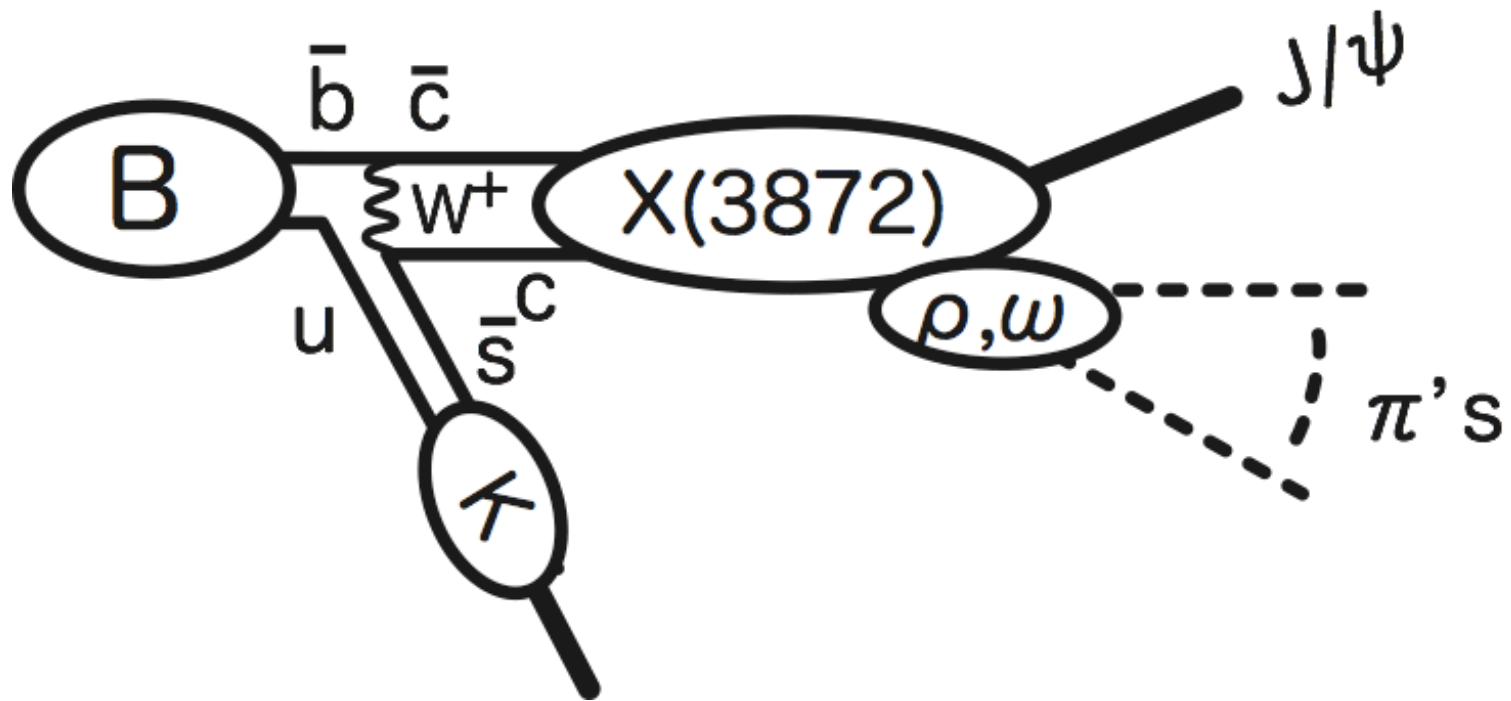
I. V. Danilkin and Y. A. Simonov, Phys. Rev. Lett., 105, 102002 (2010). **Quark model approach with $c\bar{c}$ - DD^* int.**
S. Coito, G. Rupp, and E. van Beveren, Eur. Phys. J. C, 71, 1762 (2011). **Hadronic model approach with ρ -J/psi, ω -J/psi**

Charmonium-molecule hybrid structure has following properties

- Large isospin symmetry breaking
- No charged partner
- Reasonable production rate
- Consistent with Z_b

$$B^+ \rightarrow K^+ + J/\psi + \pi\pi(\pi)$$

- $B^+ \rightarrow \underline{X(3872)} + K^+$
 $\rightarrow J/\psi + \underline{\text{vector meson}} \rightarrow \pi's$



$$\frac{Br(X(3872) \rightarrow \pi^+ \pi^- \pi^0 J / \psi)}{Br(X(3872) \rightarrow \pi^+ \pi^- J / \psi)} = 1.0 \pm 0.4 \pm 0.3$$

With kinematical suppression factor correction, Amplitude ratio is

$$\frac{A(X(3872) \rightarrow \rho J / \psi)}{A(X(3872) \rightarrow \omega J / \psi)} = 0.27 \pm 0.02 \quad \text{M. Suzuki, Phys. Rev. D72: 114013, 2005}$$

- Isovector component is smaller than isoscalar component : 25 - 30%
- Estimation of isospin component from this value is an issue of the discussion

D. Gamermann and E. Oset, Phys. Rev. D80:014003,2009.

M. Kerliner and H. J. Lipkin, arXiv:1008.0203.

K. Terasaki, Prog. Theor. Phys. 122:1205,2010.

Problems of $X(3872)$ as $C \bar{C}$ State

1. Estimated energy of $2 \ ^3P_1 \ c \bar{c}$ state by the potential model is 3950 MeV, which is about 80 MeV higher than the observed mass of $X(3872)$.
2. If $X(3872)$ is $c \bar{c}$ state, it is isoscalar.
 $X(3872) \rightarrow \rho^0 J/\psi \rightarrow \pi^+ \pi^- J/\psi$: isovector
This decay means large isospin breaking.

X(3872) as $D^0 \bar{D}^{*0}$ Molecule

- $m_{D^0} + m_{\bar{D}^{*0}} = (3871.84 \pm 0.33) \text{ MeV}$
- $m_{X(3872)} = (3871.68 \pm 0.17) \text{ MeV}$

X(3872) is a very shallow bound state of $D^0 \bar{D}^{*0}$: $D^0 \bar{D}^{*0}$ Molecule

Problem of $X(3872)$ as $D^0 \bar{D}^{*0}$ -bar Molecule

1. $D^0 \bar{D}^{*0}$ -bar is 50% isovector and 50% isoscalar: Too big the isovector component
2. Why are there no charged $X(3872)$?
 $D^+ \bar{D}^{*0}$ -bar, $D^0 \bar{D}^{*-}$ molecules
3. The production rate of such molecular-like state may be too small.

Why so large isospin symmetry breaking?

- $m_{D_0} + m_{D^*_0} = (3871.84 \pm 0.33) \text{ MeV}$

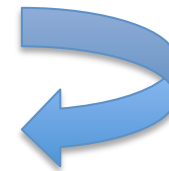
- $m_{D^+_+} + m_{D^*_{-}} = (3879.72 \pm 0.54) \text{ MeV}$

- $m_{\chi} = 3871.57 \text{ MeV}$

- Binding Energy

Neutral D case: 0.27 MeV

Charged D case: 8.15 MeV



Large
difference

Coupling between C C-bar core and $D^0 D^{*0\text{-bar}}$, $D^+ D^{*-}$

