On the importance of the tail of proton charge density

or: how to get the *rms*-radius from (e,e) data?

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Proton *rms*-radius

important quantity needed to interpret hyper-precise atomic transition energies traditionally determined via electron scattering, q = 0 slope of $G(q^2)$ analysis of world data yields $R=0.886\pm0.008\,fm$

Recent result from Lamb-shift in muonic Hydrogen

very precise radius: R=0.8418 \pm 0.0007 fm disagrees with (e,e) by many σ

Reasons for discrepancy?

many ideas discussed in literature too many to detail here no culprit identified

Investigated here:

scrutinize determination of *rms*-radius from (e,e)-data

How to determine the rms-radius?

priori this looks simple:

An unavoidable problem:

cannot measure down to q = 0even if could, finite size effect too small: $G(q) = 1 - q^2 R^2/6 + ...$ at very low q measure only the "1" given exp. uncertainties δG $\delta < r^2 >^{1/2}$ 0.004 q-region sensitive to rms-radii $0.5 < q < 1.3 fm^{-1}$ $\delta < r^4 > 1/4$ $\delta r4$ 0.002 $0.01 < Q^2 < 0.06 GeV^2/c^2$ δr2, 0.000 -0.0023.0 Õ.0 0.5 1.0 1.5 2.0 2.5 $q (fm^{-1})$

extrapolation to q = 0 introduces model dependence particularly bad for proton

Proton = particularly difficult case

form factor ~ dipole $1/(1+q^2c^2)^2$ \rightarrow density ~ exponential ~ $e^{-r/c}$

exponential density has very long tail!

Study $[\int_0^{r_{cut}}
ho(r) \; r^4 dr / \int_0^\infty
ho(r) \; r^4 dr]^{1/2}$ as function of cutoff r_{cut}



both normalized to rms-radius R

to get 98% of R must integrate out to $r \sim 3 \cdot R \sim 3 fm$

Consequence: R sensitive to $\rho(r)$ at very large r where $\rho(r)$ poorly determined affects G(q) at very low q, below q_{min}

But there are worse pitfalls!

discuss starting from two recent results:

- Inverse Polynomial fit of Bernauer *et al.*
- Continued Fraction fit of Lorenz *et al.*

Inverse Polynomial Bernauer $G(q^2) = 1/(1 + a_1q^2 + a_2q^4 +)$



Curious behavior:

between order N=7 and N=10 R^M jumps from 0.68fm to 0.96fm χ^2 best for N=10 would nominally be *the* best fit!

Bernauer *et al.* chose order N=7 ($\chi^2 \pm \text{stabilized}$)

Question remains:

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what is responsible for jump?
how can the q^{20}-term affect the rms-radius?
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Understanding



 G_M for N=10 has pole at $q > q_{max}$

In $\rho(r)_m$ this leads to oscillations extending to very large r



This affects $G_m(q^2)$ at very low q^2 , below q_{min}^2 of data



Structure at $q < q_{min}$ gives better χ^2 than N=7 (note: data are floating)

confirms old insight that absolute σ 's *much* more valuable than floating ones

Conclusion: N=10 fit is pathological. but is N=7 better?

A priori: yes, since more 'reasonable'

however: N=7 has pole too!



but pole is at larger q, happens to have much smaller effect Cannot believe either radius! Continued Fraction fits by Lorenz et al.

$$G(q) = rac{1}{1+rac{q^2b_1}{1+rac{q^2b_2}{1+\cdots}}}$$

many fits of Bernauer data with variable q_{max}

for e.g. 5 terms and $q_{max} = 3.5 fm^{-1}$ find charge-rms-radius 0.84 fm

disagrees with "accepted" result of 0.88 fm

One reason

 $\chi^2 \sim 1.4/dof$ not very good \rightarrow systematic deviations at low qSpline fit gives 1.06/dof from such a fit cannot draw conclusions



Main problem of Lorenz et al.



Extreme demonstration case

own 3-parameter Pade-fit of Bernauer data, $q < 2 fm^{-1}$ excellent $\chi^2 \sim 1.06/{\rm dof}$ no pole



rms-radius = 1.49fm!!

visible by nude eye in G(q) at very low q problem enhanced by floating data

Again: problem due to uncontrolled behavior of $G(q > q_{max})$

leads to structure of $G(q < q_{min})$ affecting q = 0 slope

For understanding: compare approach for $A>2 \quad \Longleftrightarrow \quad A\leq 2$

- for A>2 parameterize $\rho(r)$, fit to data, get *rms*-radius from integral over r
- for A ≤ 2 parameterize G(q), fit to data, get *rms*-radius from q = 0 slope

Not equivalent !!

- $\rho(r)$ automatically confined to $r < r_{max}$ by parameterization Fermi density, Gauss density, Fourier-Bessel, SOG, good physics: $\rho(r)$ must fall like $W(\kappa r)^2/r^2$ κ given by removal energy of lightest, least bound charged constituent
- constraint is missing when parameterizing G(q) $G(q > q_{max})$ can imply large ρ at large rallows for unphysical structure of G(q) below q_{min} can falsify rms-radius

This is a generic problem. MUST be avoided.

..... and unfortunately concerns most current fits

Least affected: fits including data up to maximal q of data in this case data fix large-r behavior pretty well $G(q > q_{max})$ constrained by small $G(q \sim q_{max})$ if $G \sim q^{-4}$ (\rightarrow regular $\rho(0)$)

"Solutions"

- 1. Parameterize G(q), always compute $\rho(r)$, check large-r behavior difficult as above examples show, not possible for parameterizations without FT
- 2. Parameterize $\rho(r)$ with sensible large-r fall-off $FT \rightarrow G(q)$, fit parameters to σ 's complicates life, but only a bit. Tricky point: definition of "sensible"

Solution

- parameterize $\rho(r)$ in basis with analytic FT SOG, Hermite, Laguerre, ...
- constrain ρ(r ≫) using physical model, for r where ρ(r) < 0.01 · ρ(0) fall-off of ρ given by least-bound Fock component of proton = n+π⁺ (+complications if desired, see I.S., Prog.Part.Nucl.Phys. 67 (2012)473) → adds physics explicitly, safest choice!
- fit data up to maximal q, so data constrain tail of ρ as well straightforward with above bases

To conclude

Bad news

parameterized G(q)'s may have problems very difficult to identify if this is the case or not particularly if G(q) has no FT (such as popular sum of powers of q^2) q = 0 slope could be right or wrong q = 0 slope could be sensible or not

can be believed only if $\rho(r)$ at large r has been studied! which is never dom

Good news

r-space fit with large-r constraint gives stable radii, free of diseases discussed For data=world, =world+Bernauer(+.4%), =world with fixed or floating norm, find

 $R_{ch} = 0.886 \pm .008 fm, \qquad R_m = 0.858 \pm .024 fm$

see I.S., Prog. Part. Nucl. Phys. 67 (2012) 473

.....unfortunately it does not help with μH discrepancy