

Coupling LIT and CC methods: towards continuum spectra of “not-so-few”-body systems

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Summary:

- Motivation
- Outline of LIT and CC methods and their coupling
- Validation on ^4He photodisintegration cross section
- Results on the Giant Dipole Resonance of ^{16}O and more

Work done in collaboration with

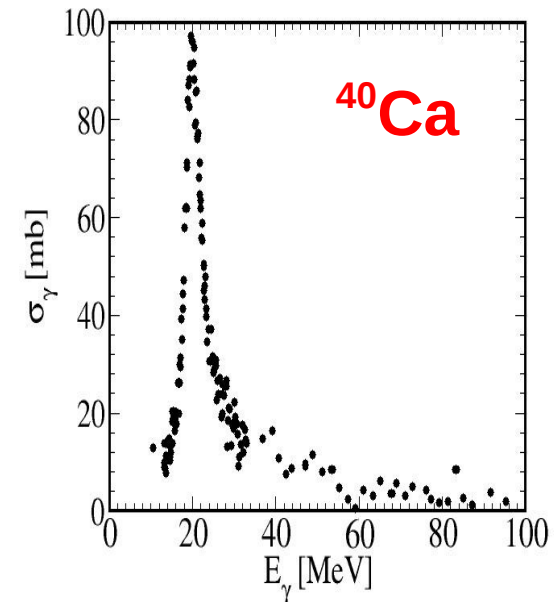
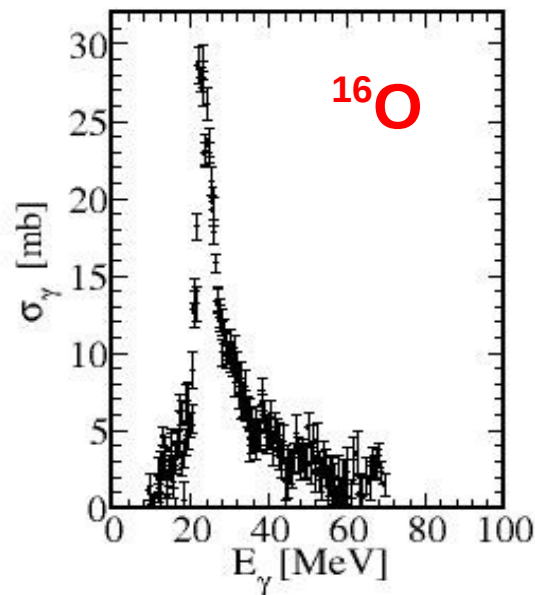
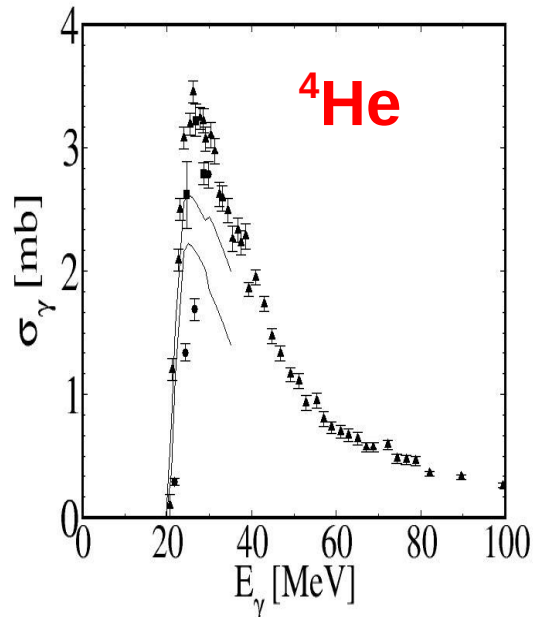
- Sonia Bacca (TRIUMF)
- Nir Barnea (Jerusalem)
- Gaute Hagen (ORNL)
- Thomas Papenbrock (ORNL)
- Mirko Miorelli (Trento)

Motivations

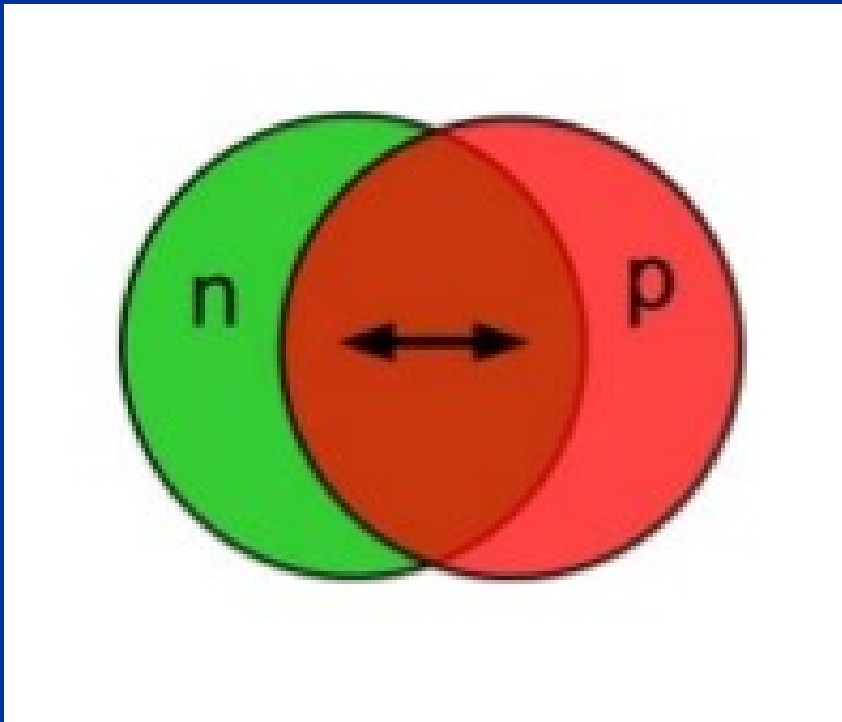
- **Continuum** spectra of “more than 3”- body systems is a great challenge for theory. On the other hand...
- ... for increasing N some spectra show **very interesting structures** which have simple physical interpretations (**collective modes** ??)
- The interesting question is:
*is an **N-N realistic potential**, as only ingredient of an **ab initio approach**, able to reproduce such structures ??*

A famous example: the so called **Giant Dipole Resonance (GDR)**

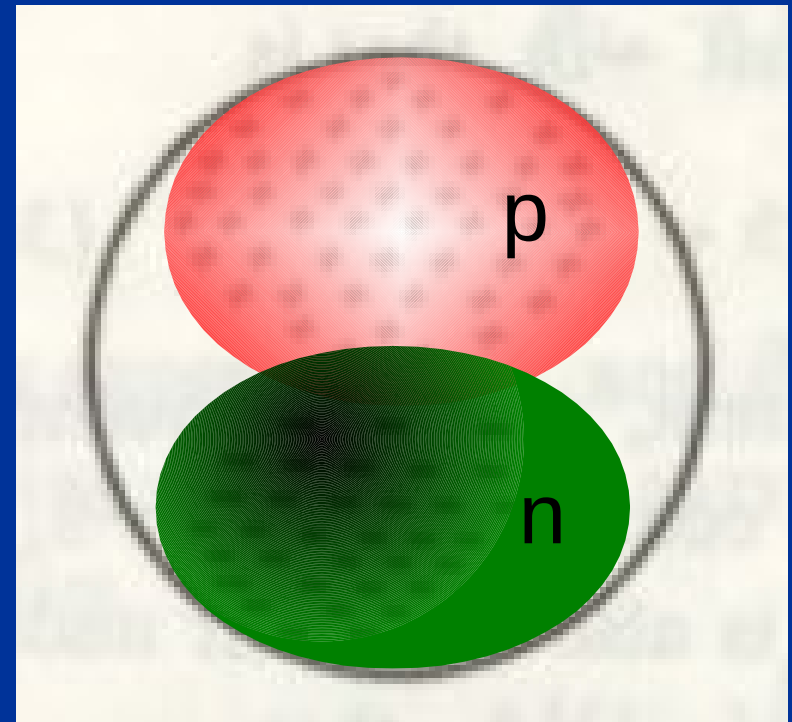
It is a pronounced **low energy** structure exhibited by the **total photoabsorption** cross section in the **continuum** all over the nuclear table !!!



Historically the **GDR** has been interpreted as
an harmonic **collective motion**
of protons against neutrons



Gamow - Teller **model** 1946



Steinwedel - Jensen **model** 1959

In the following decades various many-body theories (RPA and various improvements) based on *effective interactions* (Skyrme and similar) have tried to account for the **GDR**

The question is whether **ab initio** approaches in terms of **N-N realistic potentials** are able to reproduce it

- The **GDR** is in the continuum part of the spectrum
- The **Lorentz Integral Transform (LIT)** method reduces the **continuum** problem to a **bound state-like** problem
- Up to now the **LIT** bound-state-like equations have been formulated and solved within the **FY (A=3)**, **HH/EIHH (A=4,6,7)**, **NCSM (A=4)**
- A good ab initio method for bound state properties of larger A systems is the **Coupled-Cluster (CC)**

we can try to formulate the LIT bound-state-like equations within the CC method

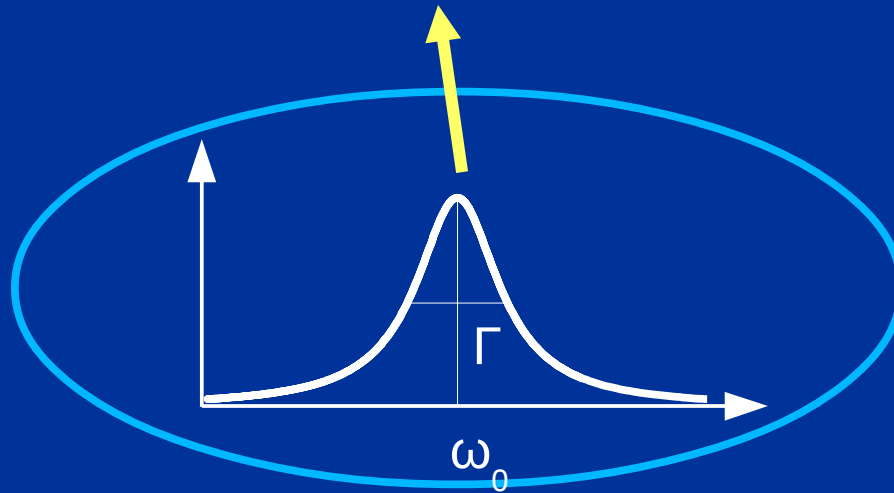
Remind the LIT method (in 3 lines):

$$\sigma(E_\gamma) = 4\pi^2 \alpha \sum_n |\langle n | D | 0 \rangle|^2 \delta(E_\gamma - E_n + E_0)$$

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Theorem:

$$\mathbf{L}(\omega_0, \Gamma) = \langle \tilde{\Psi} | \tilde{\Psi} \rangle \quad (\langle \infty)$$

where

$$[H - \omega_0 - i\Gamma] | \tilde{\Psi} \rangle = D | 0 \rangle$$

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Bound state like equation!

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$$\langle \tilde{\Psi} | \tilde{\Psi} \rangle =$$

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Similarity Transf. ANSATZ!

$|SD\rangle$

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$$0 = \langle 2p-2h | E |SD\rangle = \langle 2p-2h | \overline{H} |SD\rangle$$

⋮
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$$0 = \langle kp-kh | E |SD\rangle = \langle kp-kh | \overline{H} |SD\rangle$$

CC equations

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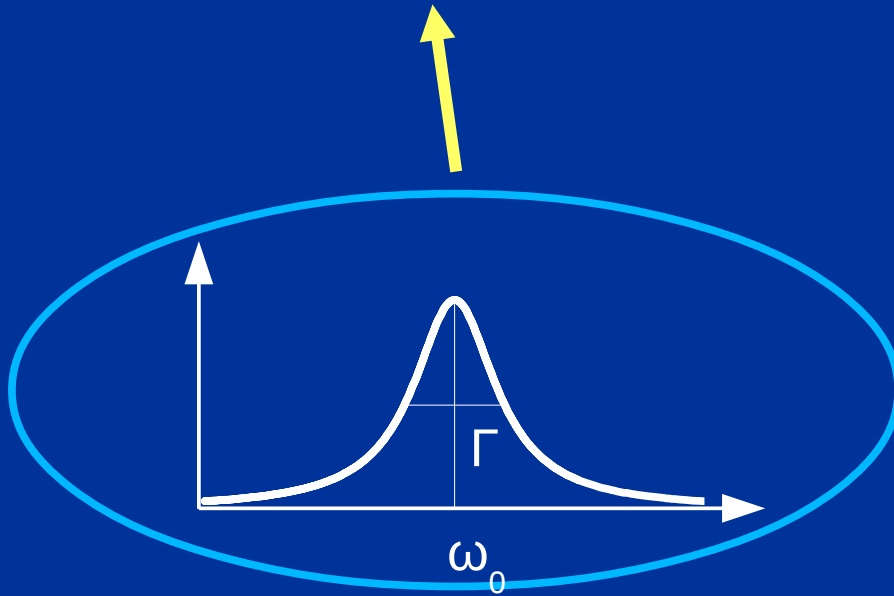
$$0 = \langle 2p-2h | E |SD\rangle = \langle 2p-2h | \overline{H} |SD\rangle$$

$$0 = \langle 3p-3h | E |SD\rangle = \langle 3p-3h | \overline{H} |SD\rangle$$

CCSDT

LIT+CC:

$$\mathbf{L}(\omega_0, \Gamma) = \int dE_\gamma [(E_\gamma - \omega_0)^2 + \Gamma^2]^{-1} \sum_n |\langle n | \mathbf{D} | 0 \rangle|^2 \delta(E_\gamma - E_n + E_0)$$



Writing the Lorentzian kernel $[(E_\gamma - \omega_0)^2 + \Gamma^2]^{-1}$ as a product

$$[H - E_0 - \omega_0 - i\Gamma]^{-1} * [H - E_0 - \omega_0 + i\Gamma]^{-1}$$

And using the delta function in the integral and completeness of eigenstates ($|n\rangle\langle n| = 1$):

$$\mathbf{L}(\omega_0, \Gamma) = \langle 0 | \mathbf{D} [H - E_0 - \omega_0 - i\Gamma]^{-1} [H - E_0 - \omega_0 + i\Gamma]^{-1} \mathbf{D} | 0 \rangle$$

$$= \langle \tilde{\Psi} | \tilde{\Psi} \rangle$$

Now use the similarity transformations inserting
 $e^T e^{-T} = 1$

$$\mathbf{L}(\omega_0, \Gamma) =$$

$$= \langle 0 | \underset{\uparrow}{D} \underset{\uparrow}{[H - E_0 - \omega_0 - i\Gamma]^{-1}} \underset{\uparrow}{[H - E_0 - \omega_0 + i\Gamma]^{-1}} \underset{\uparrow}{D} | 0 \rangle$$

Insert $e^T e^{-T} = 1$

$$\mathbf{L}(\omega_0, \Gamma) =$$

$$= \langle 0 | \mathbf{e}^T \overline{\mathbf{D}} [\overline{\mathbf{H}} - E_0 - \omega_0 - i\Gamma]^{-1} [\overline{\mathbf{H}} - E_0 - \omega_0 + i\Gamma]^{-1} \overline{\mathbf{D}} \mathbf{e}^{-T} | 0 \rangle$$

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$$\mathbf{L}(\omega_0, \Gamma) =$$

$$= \langle 0 | e^{\mathbf{T}} \overline{\mathbf{D}} [\overline{\mathbf{H}} - E_0 - \omega_0 - i\Gamma]^{-1} [\overline{\mathbf{H}} - E_0 - \omega_0 + i\Gamma]^{-1} \overline{\mathbf{D}} | \text{SD} \rangle$$

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$\langle 0 | e^{\mathbf{T}} \neq \langle \text{SD} | !!$

$\overline{\mathbf{H}}$ and $\overline{\mathbf{D}}$ are non hermitian operators and
 one has $| \text{SDR} \rangle$ and $\langle \text{SDL} |$

$$\mathbf{L}(\omega_0, \Gamma) = \langle \tilde{\Phi} | \tilde{\Psi} \rangle$$

$$= \langle 0 | e^{\mathbf{T}} \overline{\mathbf{D}} [\overline{\mathbf{H}} - E_0 - \omega_0 - \imath \Gamma]^{-1} [\overline{\mathbf{H}} - E_0 - \omega_0 + \imath \Gamma]^{-1} \overline{\mathbf{D}} | \text{SD} \rangle$$

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one has $| \text{SDR} \rangle$ and $\langle \text{SDL} |$

One can also avoid the second equation by:

1) rewriting the Lorentz kernel as a difference (*instead of a product*)

$$[(E_\gamma - \omega_0)^2 + \Gamma^2]^{-1} = (2\Gamma)^{-1} \{ (E_\gamma - \omega_0 - i\Gamma)^{-1} - (E_\gamma - \omega_0 + i\Gamma)^{-1} \}$$

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$$\mathbf{L}(\omega_0, \Gamma) = \langle \text{SDL} | \overline{\mathbf{D}} [\overline{\mathbf{H}} - E_0 - \omega_0 + i\Gamma]^{-1} \overline{\mathbf{D}} | \text{SDR} \rangle - \\ \langle \text{SDL} | \overline{\mathbf{D}} [\overline{\mathbf{H}} - E_0 - \omega_0 - i\Gamma]^{-1} \overline{\mathbf{D}} | \text{SDR} \rangle$$

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2) use the **non** hermitian **Lanczos** algorithm to rewrite as a continuous fraction

results

We have applied the CC-LIT method to ^{16}O GDR using $\mathbf{D} = \sum_i z_i \boldsymbol{\tau}_i^3$ and NN forces from EFT (N3LO)

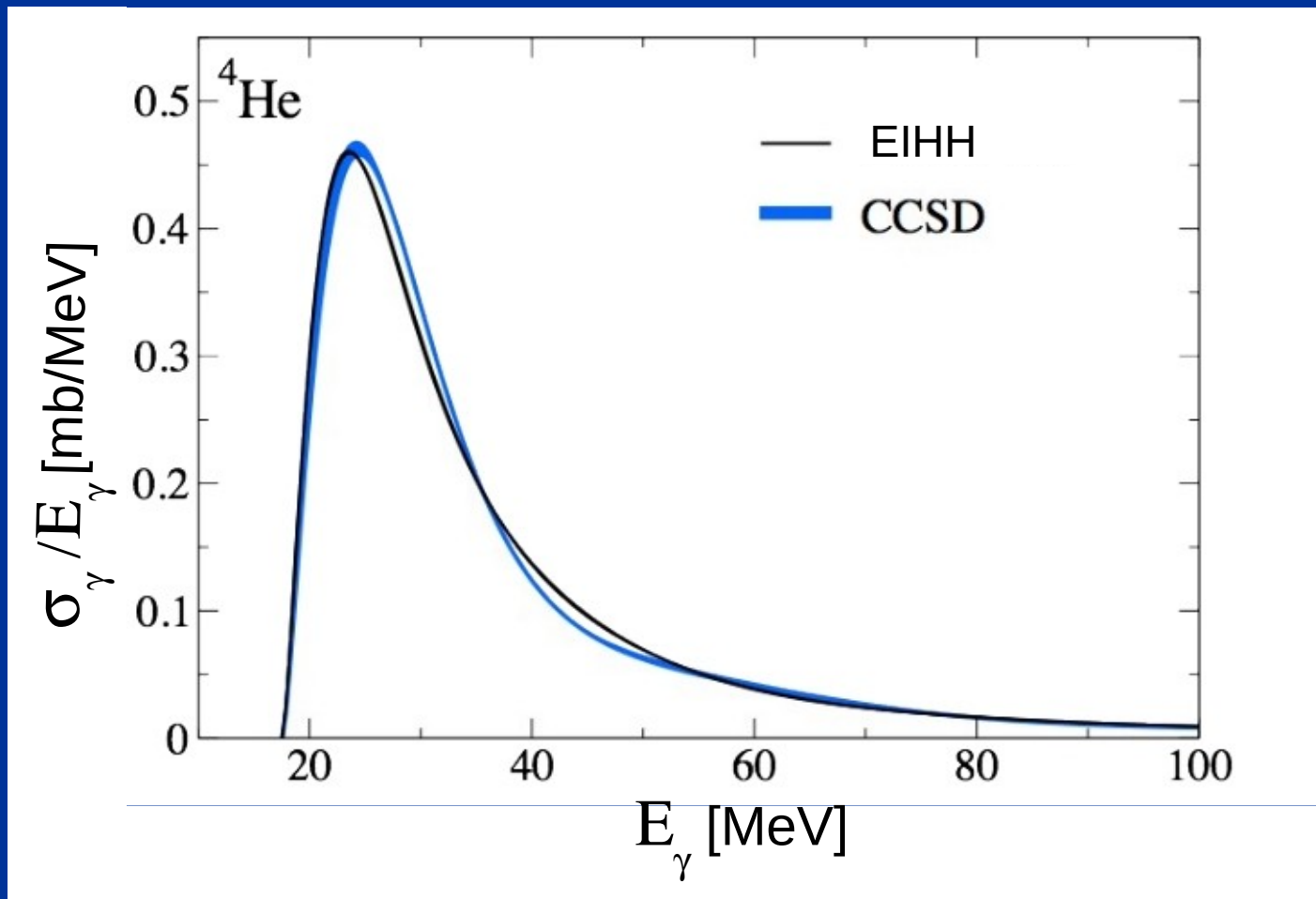
S. Bacca, N. Barnea, G. Hagen, G.O., T. Papenbrock

First principles description of the giant dipole resonance in ^{16}O

arXiv:1303.7446 [nucl-th] to appear on PRL

First we have validated the method on ${}^4\text{He}$

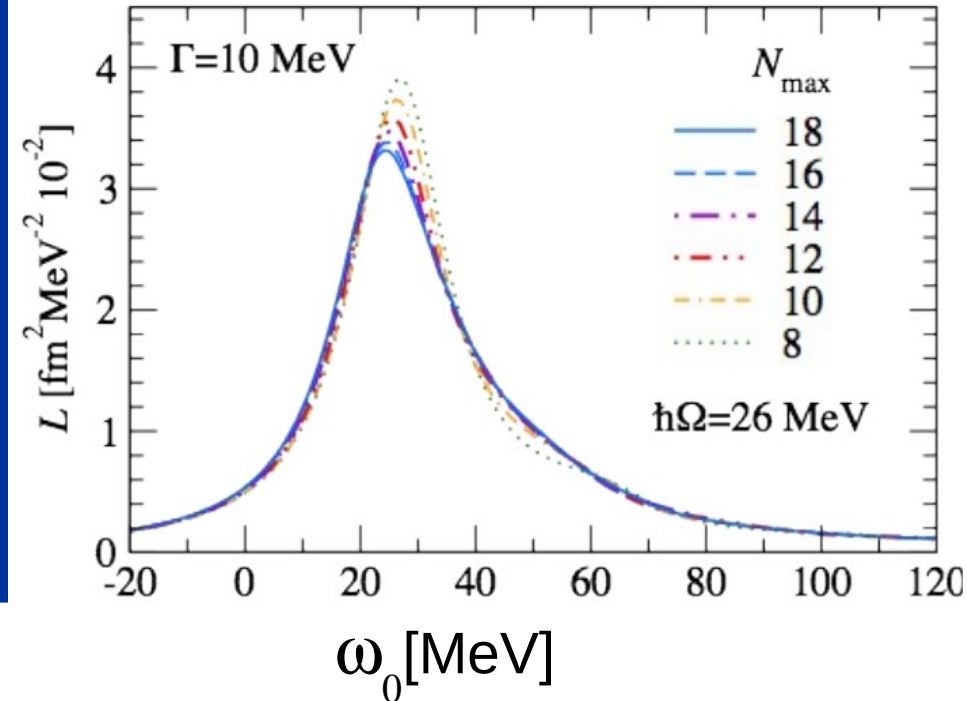
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Good agreement !

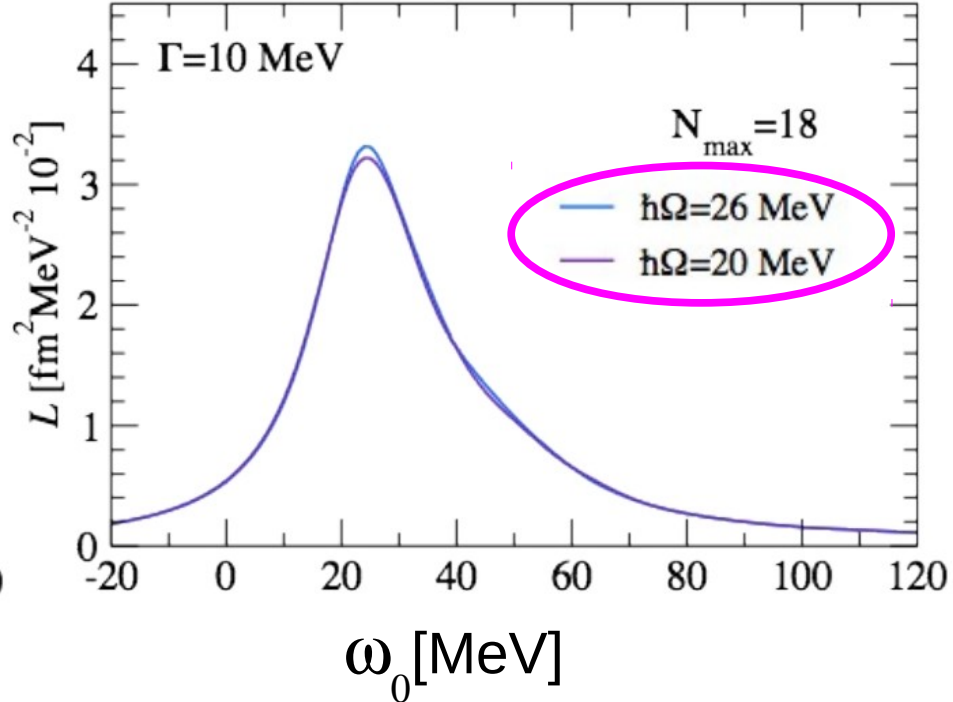
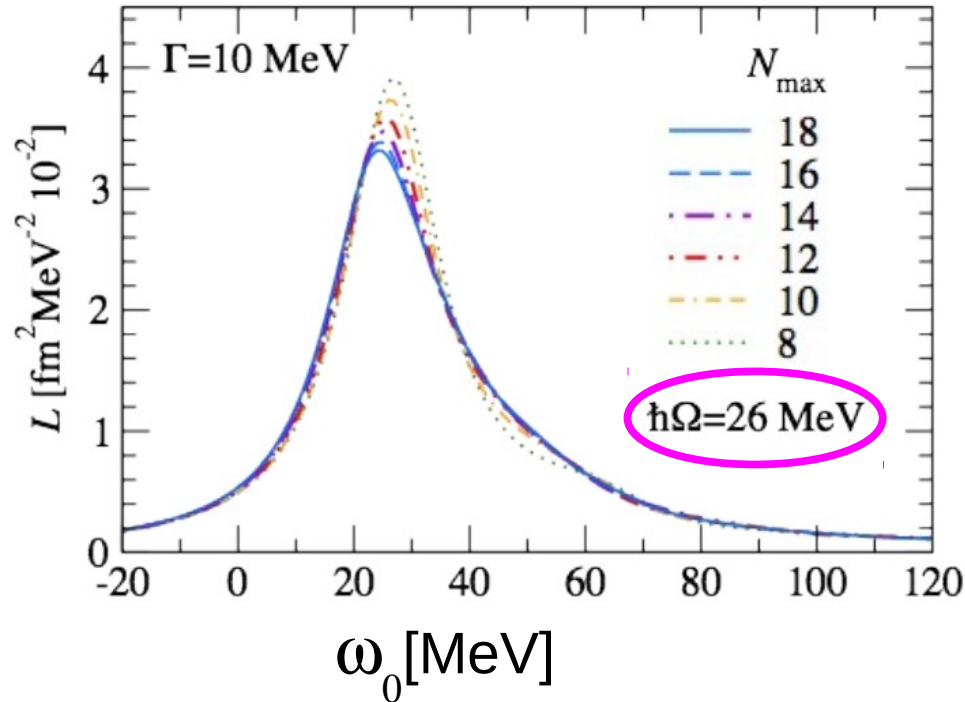
^{16}O

^{16}O : Convergence of the LIT



Good convergence!

^{16}O : H.O. parameter dependence

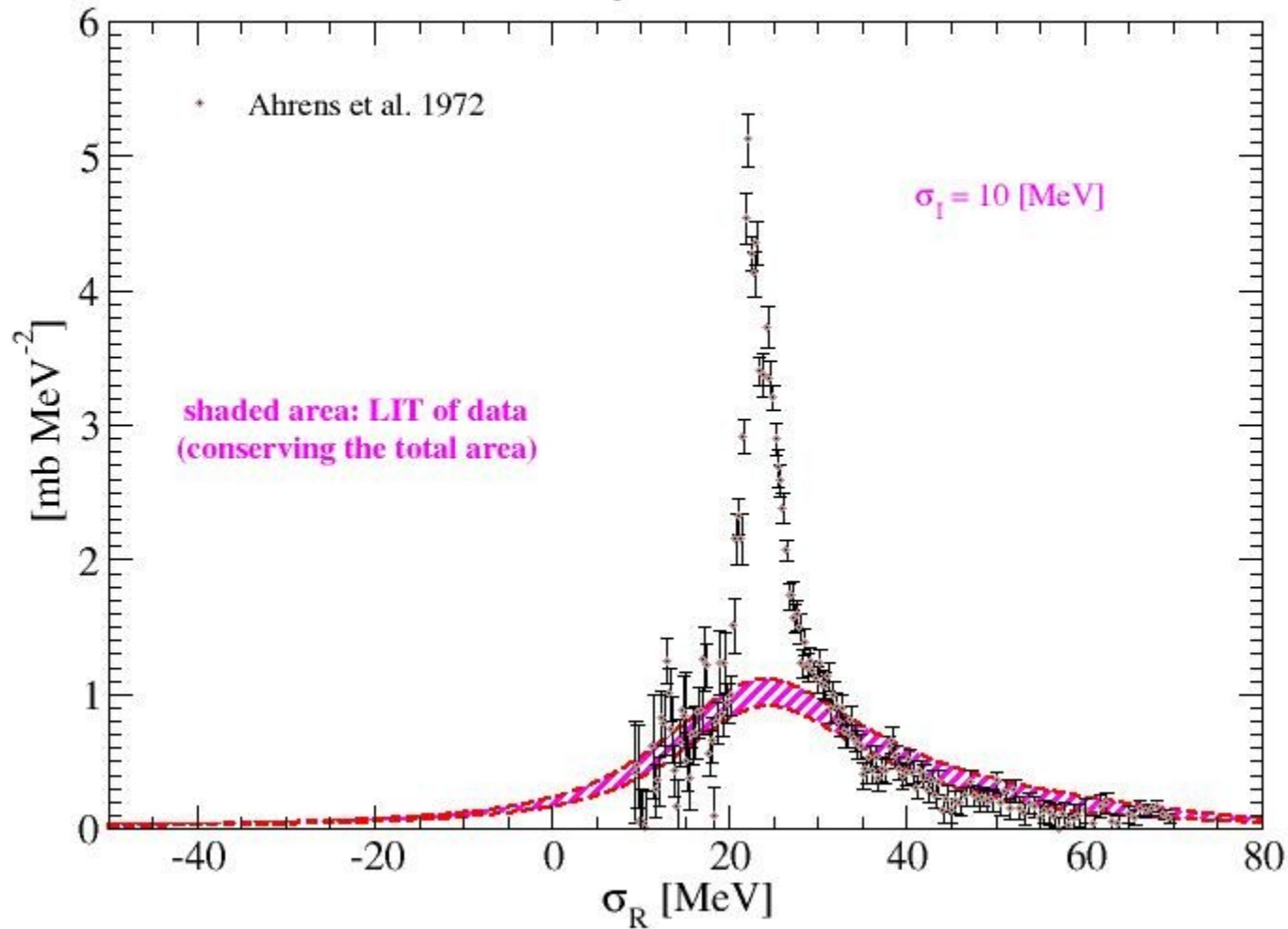


Good convergence!

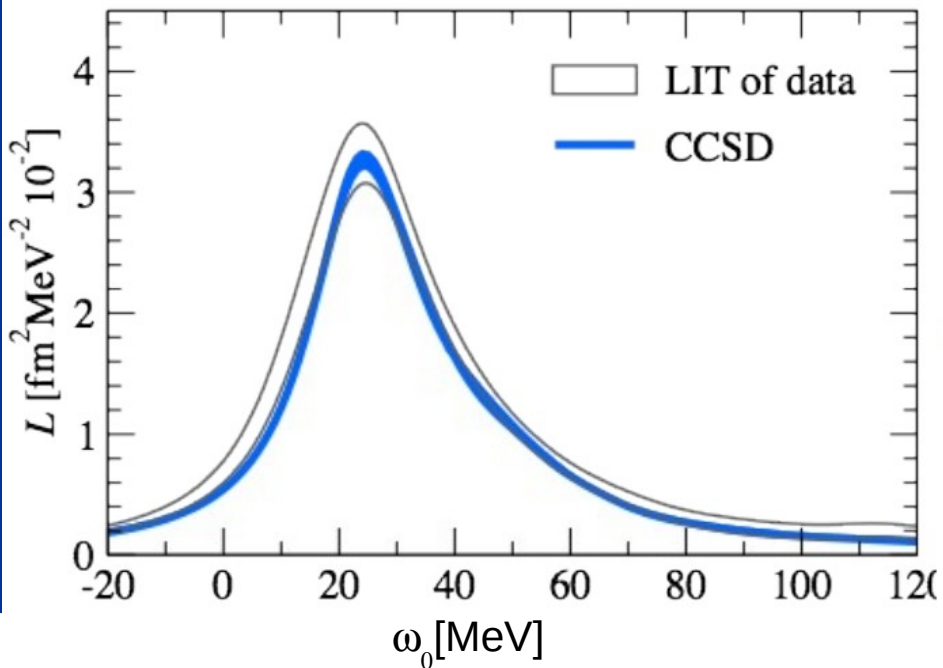
**Small HO dependence:
use it as error bar!**

Exper. LIT of the photoabsorption cross section of ^{16}O

$\sigma_I = 10$ [MeV]



Comparison between the Exp. and Theor. LIT's

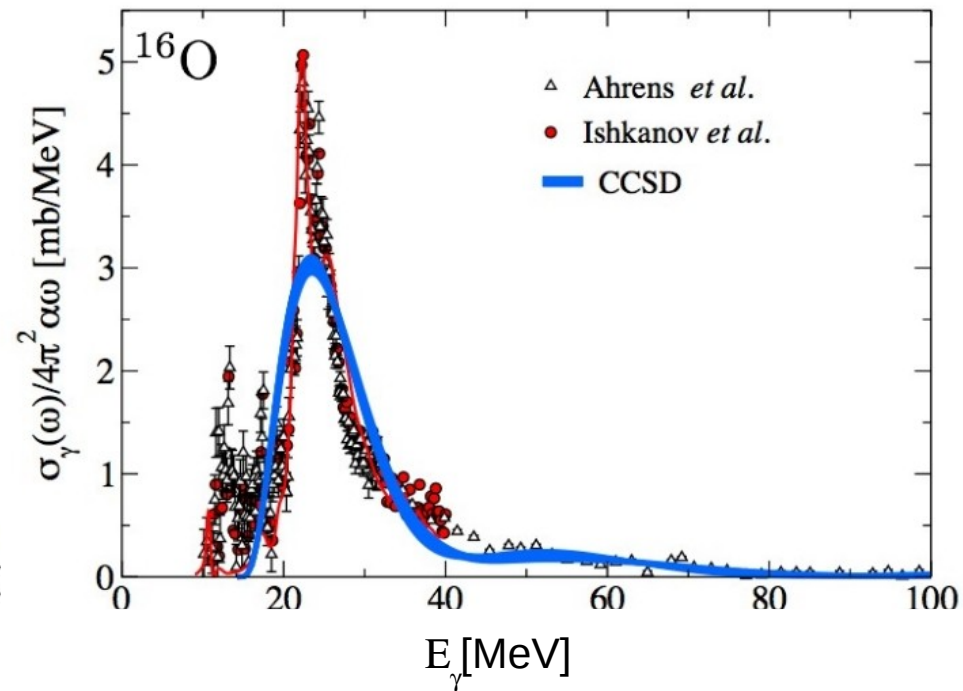
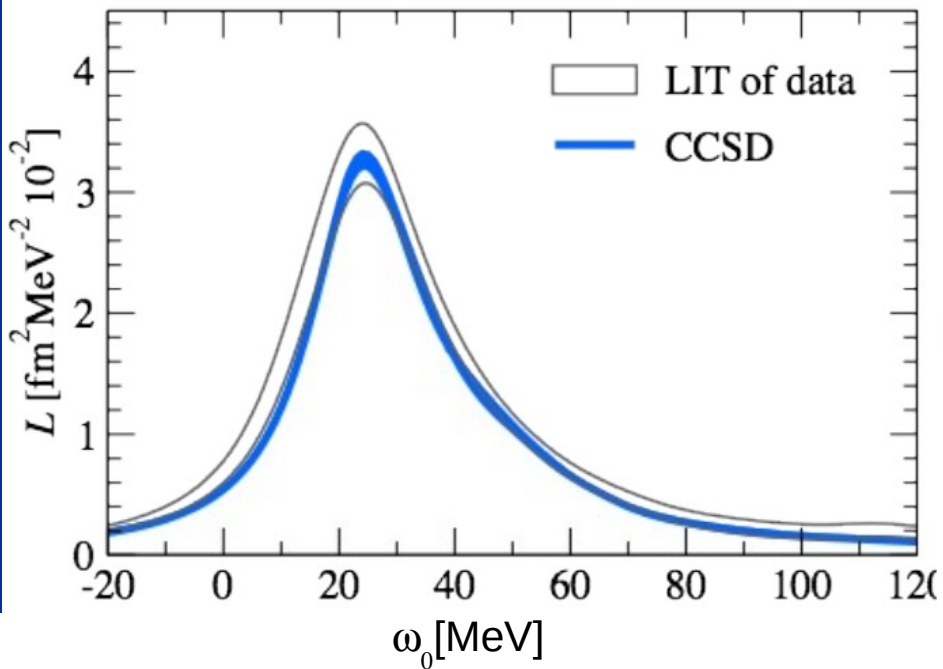


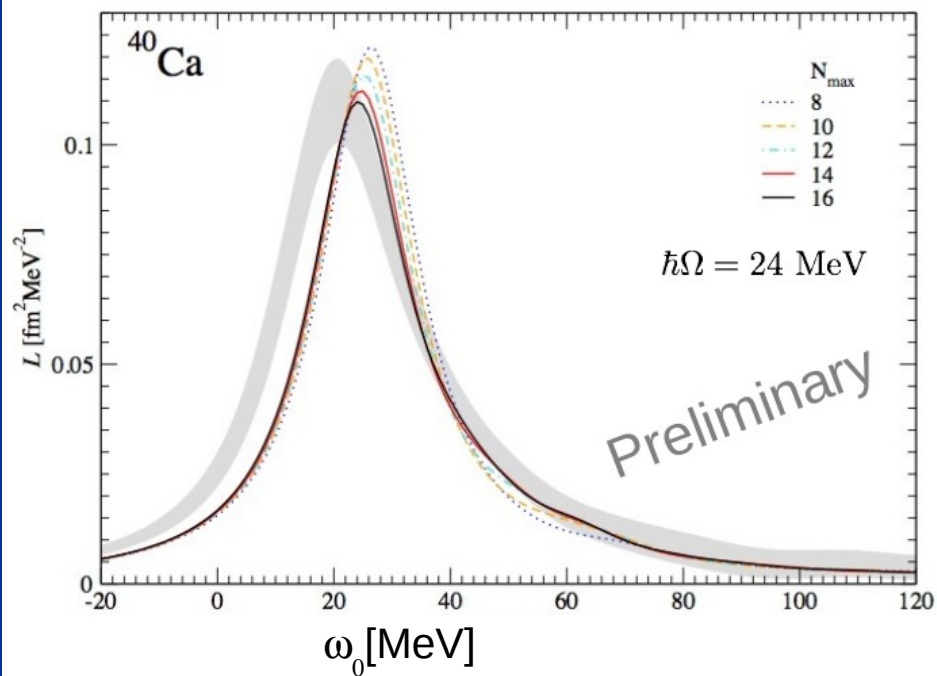
The width of the GDR is around 6 MeV

Our width is 10 MeV

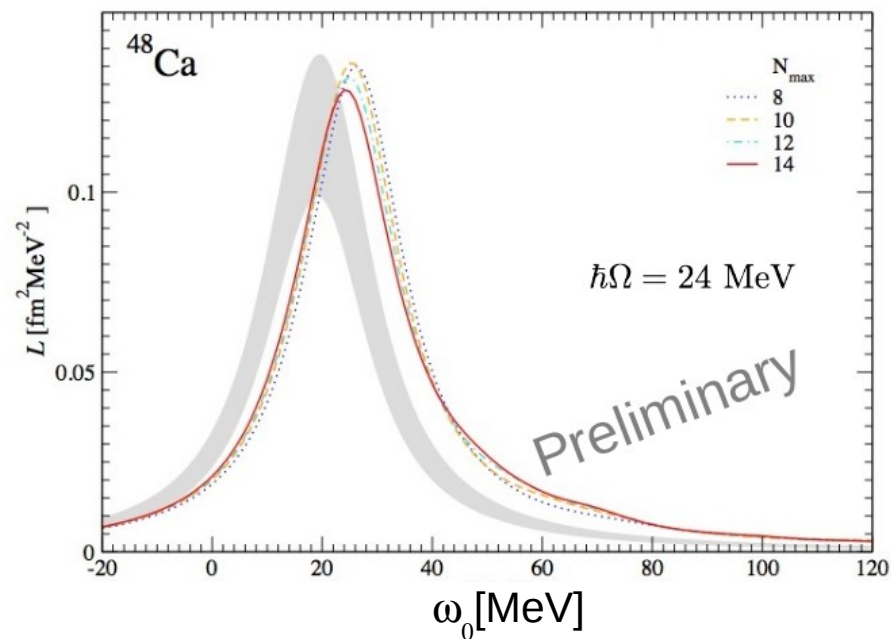
We can try to invert the LIT

Comparison between the Exp. and Theor. LIT's





Mirko Miorelli



Conclusions

- The coupling of LIT and CC methods is a viable way towards ab initio explorations of collective phenomena like “giant resonances”
- The position of the GDR of ^{16}O is reproduced correctly by a *realistic* NN potential

Perspectives

- . Study effects of NNN
- Other *Giant* resonances (monopole, quadrupole...)
- Other nuclei
-

**A lot of work to build a bridge
between
few-body and “many-body” physics**

