

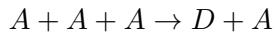
# Three-body recombination in cold gases

D.V. Fedorov,  
P.K. Sørensen, A.S. Jensen, N.T. Zinner

Aarhus Universitet

EFB22, Krakow, September 2013

# Recombination of cold atoms in traps



## Effective range expansion; Feshbach resonance

$$k \cot \delta = -\frac{1}{a} + \frac{1}{2}Rk^2 + O(k^4)$$

## Effective range expansion; Feshbach resonance

$$k \cot \delta = -\frac{1}{a} + \frac{1}{2}Rk^2 + O(k^4)$$

$$U(\mathbf{r}) = a \frac{4\pi\hbar^2}{m} \delta(\mathbf{r})$$

## Effective range expansion; Feshbach resonance

$$k \cot \delta = -\frac{1}{a} + \frac{1}{2}Rk^2 + O(k^4)$$

$$U(\mathbf{r}) = a \frac{4\pi\hbar^2}{m} \delta(\mathbf{r})$$

$$\left. \frac{1}{\psi} \frac{\partial \psi}{\partial r} \right|_{r=0} = -\frac{1}{a}$$

## Effective range expansion; Feshbach resonance

$$k \cot \delta = -\frac{1}{a} + \frac{1}{2}Rk^2 + O(k^4)$$

$$U(\mathbf{r}) = a \frac{4\pi\hbar^2}{m} \delta(\mathbf{r})$$

$$\left. \frac{1}{\psi} \frac{\partial \psi}{\partial r} \right|_{r=0} = -\frac{1}{a}$$

Using Feshbach resonances one can tune  $a$  at wish:

$$-\infty < a < \infty$$

## Effective range expansion; Feshbach resonance

$$k \cot \delta = -\frac{1}{a} + \frac{1}{2}Rk^2 + O(k^4)$$

$$U(\mathbf{r}) = a \frac{4\pi\hbar^2}{m} \delta(\mathbf{r})$$

$$\left. \frac{1}{\psi} \frac{\partial \psi}{\partial r} \right|_{r=0} = -\frac{1}{a}$$

Using Feshbach resonances one can tune  $a$  at wish:

$$-\infty < a < \infty$$

Two-channel zero-range model describes Feshbach resonances quite well.

## Recombination into shallow dimer

$$-\dot{N} = 3 \frac{N^3}{6} \frac{2\pi}{\hbar} |T_{fi}|^2 \frac{d\nu_f}{dE_f} \Big|_{E_f = \frac{\hbar^2}{ma^2}}$$



## Recombination into shallow dimer

$$-\dot{N} = 3 \frac{N^3}{6} \frac{2\pi}{\hbar} |T_{fi}|^2 \frac{d\nu_f}{dE_f} \Big|_{E_f = \frac{\hbar^2}{ma^2}}$$

$$d\nu_f = \frac{V d^3 q_f}{(2\pi)^3} \Big|_{q_f = \frac{2}{\sqrt{3}} a^{-1}} = \frac{2}{3\sqrt{3}\pi^2} \frac{Vm}{\hbar^2 a} dE_f$$

## Recombination into shallow dimer

$$-\dot{N} = 3 \frac{N^3}{6} \frac{2\pi}{\hbar} |T_{fi}|^2 \frac{d\nu_f}{dE_f} \Big|_{E_f = \frac{\hbar^2}{ma^2}}$$

$$d\nu_f = \frac{V d^3 q_f}{(2\pi)^3} \Big|_{q_f = \frac{2}{\sqrt{3}} a^{-1}} = \frac{2}{3\sqrt{3}\pi^2} \frac{Vm}{\hbar^2 a} dE_f$$

$$T_{fi} = \int d^3 r d^3 R \left[ \psi_D(r) \frac{e^{i\mathbf{q}_f \mathbf{R}}}{\sqrt{V}} \right] U(\mathbf{R} - \frac{1}{2} \mathbf{r}) \left[ \frac{e^{i\mathbf{k} \mathbf{r}}}{\sqrt{V}} \frac{e^{i\mathbf{q} \mathbf{R}}}{\sqrt{V}} \right]$$

## Recombination into shallow dimer

$$-\dot{N} = 3 \frac{N^3}{6} \frac{2\pi}{\hbar} |T_{fi}|^2 \frac{d\nu_f}{dE_f} \Big|_{E_f = \frac{\hbar^2}{ma^2}}$$

$$d\nu_f = \frac{V d^3 q_f}{(2\pi)^3} \Big|_{q_f = \frac{2}{\sqrt{3}} a^{-1}} = \frac{2}{3\sqrt{3}\pi^2} \frac{Vm}{\hbar^2 a} dE_f$$

$$T_{fi} = \int d^3 r d^3 R \left[ \psi_D(r) \frac{e^{i\mathbf{q}_f \mathbf{R}}}{\sqrt{V}} \right] U(\mathbf{R} - \frac{1}{2} \mathbf{r}) \left[ \frac{e^{i\mathbf{k} \mathbf{r}}}{\sqrt{V}} \frac{e^{i\mathbf{q} \mathbf{R}}}{\sqrt{V}} \right]$$

$$\psi_D(r) \propto a^{-1/2} \frac{e^{-\frac{r}{a}}}{r}, \quad T_{fi} \propto \frac{\hbar^2 a^{5/2}}{V^{3/2} m}$$

## Recombination into shallow dimer

$$-\dot{N} = 3 \frac{N^3}{6} \frac{2\pi}{\hbar} |T_{fi}|^2 \frac{d\nu_f}{dE_f} \Big|_{E_f = \frac{\hbar^2}{ma^2}}$$

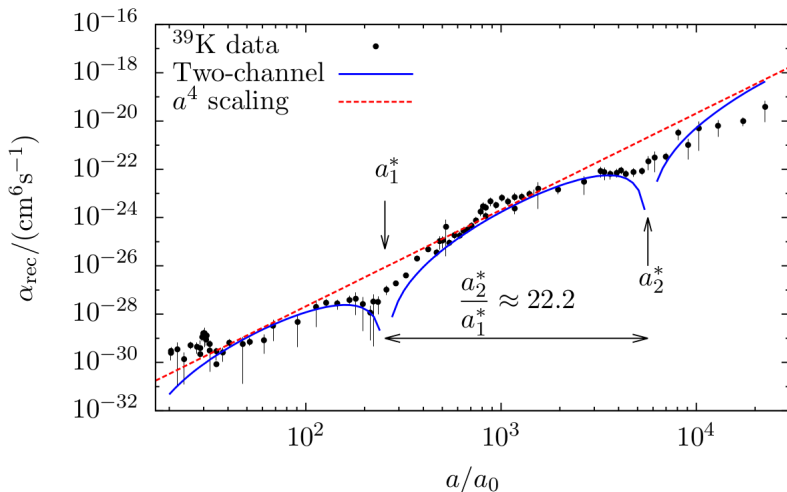
$$d\nu_f = \frac{V d^3 q_f}{(2\pi)^3} \Big|_{q_f = \frac{2}{\sqrt{3}} a^{-1}} = \frac{2}{3\sqrt{3}\pi^2} \frac{Vm}{\hbar^2 a} dE_f$$

$$T_{fi} = \int d^3 r d^3 R \left[ \psi_D(r) \frac{e^{i\mathbf{q}_f \mathbf{R}}}{\sqrt{V}} \right] U(\mathbf{R} - \frac{1}{2} \mathbf{r}) \left[ \frac{e^{i\mathbf{k} \mathbf{r}}}{\sqrt{V}} \frac{e^{i\mathbf{q} \mathbf{R}}}{\sqrt{V}} \right]$$

$$\psi_D(r) \propto a^{-1/2} \frac{e^{-\frac{r}{a}}}{r}, \quad T_{fi} \propto \frac{\hbar^2 a^{5/2}}{V^{3/2} m}$$

$$n = \frac{N}{V} \Rightarrow \boxed{\dot{n} = -\alpha_{\text{rec}} n^3}, \quad \alpha_{\text{rec}} \propto \frac{\hbar a^4}{m}$$

# Recombination into shallow dimer



**Figure:** The recombination coefficient  $\alpha_{\text{rec}}$  for  $^{39}\text{K}$  as function of the scattering length  $a$ ; experimental data from [Zaccanti 2009].

## Recombination into deep dimers

There is no more  $\psi_D(r) = a^{-1/2} \frac{e^{-\frac{r}{a}}}{r}$  but rather a number of short-range dimers.

## Recombination into deep dimers

There is no more  $\psi_D(r) = a^{-1/2} \frac{e^{-\frac{r}{a}}}{r}$  but rather a number of short-range dimers.

We shall try to describe recombination into those deep dimers using a short-range optical potential (square well, actually),

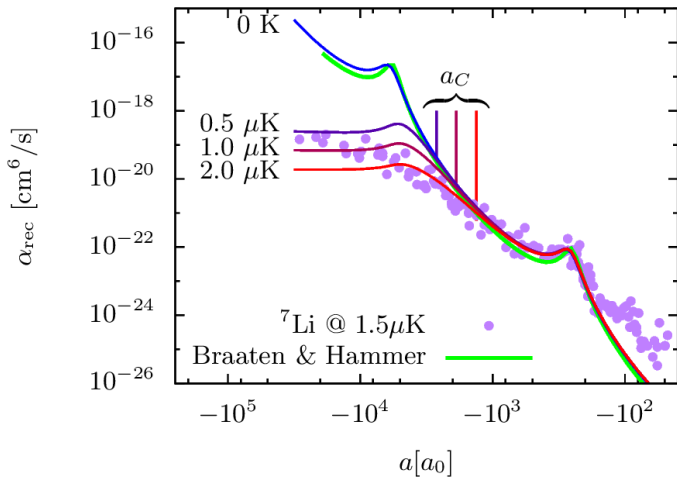
$$(1 + i\epsilon)V_{\text{short range}}(\rho), \quad \rho^2 = r_1^2 + r_2^2 + r_3^2$$

## Temperature dependence

$$\alpha_{\text{rec}}(a; T) = \frac{1}{2T^3} \int E^2 e^{-\frac{E}{T}} \alpha_{\text{rec}}(a, E) dE$$



# Recombination into deep dimers



**Figure:** The recombination coefficient  $\alpha_{\text{rec}}$  for  $^7\text{Li}$  as function of the scattering length  $a$ ; experimental data at  $1.5 \mu\text{K}$  from [Dyke 2013].