

# ${}^6_{\Lambda}\text{H}$ Modeled as ${}^4_{\Lambda}\text{H} + n + n$

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Evidence for the existence of a bound state of the hypernucleus  ${}^6_{\Lambda}\text{H}$  was recently reported by Agnello *et al.* [1]. The possible existence of such a neutron-rich hypernucleus was originally discussed by Dalitz and Levi-Setti [2]. The unanswered question was whether the  $\Lambda N$  interaction is sufficiently strong so as to bind  ${}^6_{\Lambda}\text{H}$  ( ${}^5\text{H}$  being unbound), given that the hypertriton with a deuteron core is barely bound. Dalitz and Levi-Setti suggested that  ${}^6_{\Lambda}\text{H}$  be considered as a  ${}^4_{\Lambda}\text{H} + n + n$  three-body system. Starting from the then quoted  $\Lambda$ -separation energy of 2.4 MeV for  ${}^4_{\Lambda}\text{H}$ , they added 2 times the estimated 0.9 MeV per p-shell neutron binding in  ${}^7_{\Lambda}\text{Be}$  to obtain a  $B_{\Lambda}({}^6_{\Lambda}\text{H})$  of  $\approx 4.2$  MeV. The  ${}^6_{\Lambda}\text{H}$   $\Lambda$ -separation energy ( $\rightarrow {}^5\text{H} + \Lambda$ ) estimate from Ref. [1], is  $4.0 \pm 1.1$  MeV. This lies close to the value of the  ${}^6_{\Lambda}\text{He}$   $\Lambda$ -separation energy ( $\rightarrow {}^5\text{He} + \Lambda$ ) of  $B_{\Lambda}({}^6_{\Lambda}\text{He}) = 4.18 \pm 0.10$  MeV [3], even though the isospin of the former is 3/2 and the isospin of the latter is 1/2. In fact, if one takes the  ${}^4_{\Lambda}\text{H}$   $\Lambda$ -separation energy to be the  $2.04 \pm 0.04$  MeV quoted in Ref. [3] and combines it with the two-neutron energy difference between the  ${}^7_{\Lambda}\text{He}$  and  ${}^5_{\Lambda}\text{He}$  isotopes, one obtains an estimate for the  $\Lambda$ -separation energy of  ${}^6_{\Lambda}\text{H}$  of  $\sim 4.02 \pm 0.4$  MeV. To arrive at this result, one must note that Davis clearly states that the  $\Lambda$ -separation energy cannot be obtained by averaging the distinctly different observed values [4]. Pniewski and Danysz point out that one set of measurements clusters around an average of  $5.1 \pm 0.4$  MeV corresponding to the ground state, while a second set of measurements clusters around  $3.2 \pm 0.4$  MeV, corresponding to the first excited state of the core  ${}^6\text{He}$ . Thus, the difference between  $B_{\Lambda}({}^5_{\Lambda}\text{He}) = 3.12 \pm 0.02$  MeV and  $B_{\Lambda}({}^7_{\Lambda}\text{He}) = 5.1 \pm 0.4$  MeV suggests that the contribution of the two p-shell neutrons to the ground-state and excited state  $\Lambda$ -separation energies of  ${}^6_{\Lambda}\text{H}$  should be about 2 MeV. Therefore, an estimate of  $B_{\Lambda}({}^6_{\Lambda}\text{H}) = 4.04 \pm 0.04$  MeV is attained, which is close to the 4.2 MeV from Ref. [2] and agrees with the experimental value of  $4.0 \pm 1.1$  MeV from Ref. [1]. A similar theoretical analysis is reported in Ref. [5]; this analysis is based on extrapolated values for  $B_{\Lambda}({}^7_{\Lambda}\text{He})$  and yields a value for  $B_{\Lambda}({}^6_{\Lambda}\text{H}) \approx 4.28$  MeV. Nevertheless, it should be noted that a different value for  $B_{\Lambda}({}^7_{\Lambda}\text{He})$  of  $5.68 \pm 0.03 \pm 0.25$  MeV obtained from a JLab experiment has been recently published [6]. This result would suggest a larger nominal value for  $B_{\Lambda}({}^6_{\Lambda}\text{H})$ , but one that is still consistent with the reported experimental value.

In addition to this interesting analysis of the possible ground state of  ${}^6_{\Lambda}\text{H}$  is the observation [1] that the mass estimates from the production and decay analysis appear to differ by about 1 MeV, which is very similar to the difference in the binding energies of the  ${}^4_{\Lambda}\text{H}$   $0^+$  ground state and  $1^+$  excited state. That is, it is postulated [1] that the  ${}^6_{\Lambda}\text{H}$  ( $1^+$ ) state is originally produced in the  $(K^-_{stop}, \pi^+)$  reaction, and this is followed by  $\gamma$  decay to the  ${}^6_{\Lambda}\text{H}$  ( $0^+$ ) ground state, whose weak decay is then observed.

Motivated by the recent experiments, we explore a three-body calculation for the  ${}^6_{\Lambda}\text{He}$  and  ${}^6_{\Lambda}\text{H}$  hypernuclei in which the core is  ${}^4_{\Lambda}\text{H}$ . Two of the interactions in the  ${}^6_{\Lambda}\text{He}$  system

$[\Lambda^4\text{H} + p + n]$  are, in principle, known. That is, the  $np$  interaction is known and the  $\Lambda^4\text{H} - p$  interaction can be fitted to the  ${}^5_\Lambda\text{He}$  binding energy. The  $\Lambda^4\text{H} - n$  interaction is to be determined by fitting to the  ${}^6_\Lambda\text{He}$  binding energy. Given the paucity of data to constrain our  $\Lambda^4\text{H} - n$  interaction, we choose to model it in a manner similar to that for the  ${}^4\text{He} - n$  interaction. In the latter case the s-wave potential is modeled alternatively as a repulsive potential [7] or an attractive potential with a forbidden bound state [8]. We examine the effect of this alternative on a  ${}^4\text{He} + n + \Lambda$  model of  ${}^6_\Lambda\text{He}$ , because it comes into play also in our  $\Lambda^4\text{H} + n + n$  model of  ${}^6_\Lambda\text{H}$  where the p-shell neutrons are Pauli blocked from the s-shell of the  $\Lambda^4\text{H}$  core.

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